

EXPERIMENTAL INVESTIGATION OF THE ONSET OF CAVITATION IN A QUADRANT-EDGE ORIFICE METER. The onset of cavitation was determined by the graphical method suggested. The validity of the method suggested.

# A GRAPHICAL METHOD TO LOCATE THE ONSET OF CAVITATION

By

ANALYSIS OF THE PROBLEM. M. V. Ramamoorthy,

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Referring to Figure 1, the flow of the fluid is

$P_1$  and  $V_1$  completely Bangalore 12, India

and  $V_2$  shows up in the discharge coefficient  $C$  and so in pipe

Reynolds' number  $R$ . K. Seetharamaiah,

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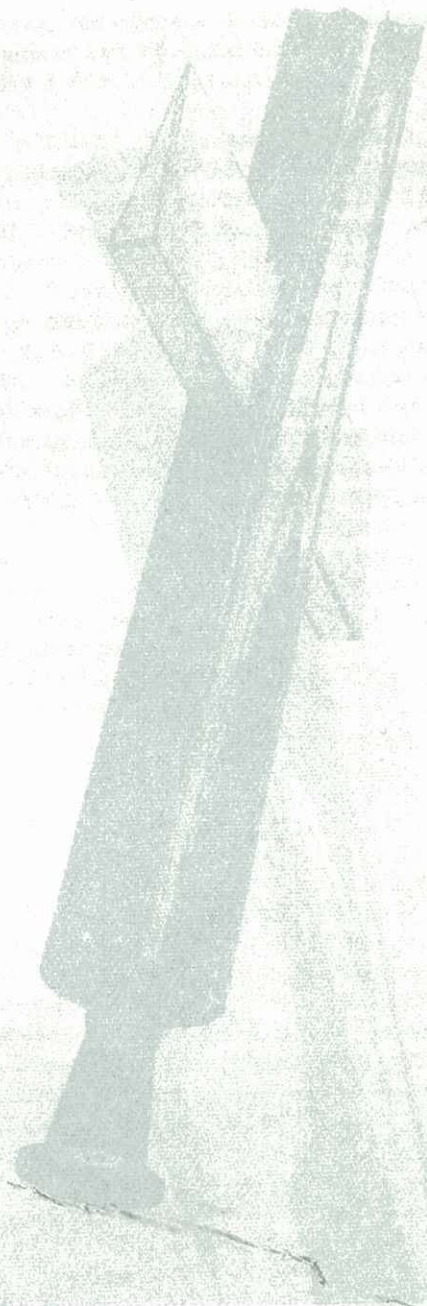
## ABSTRACT

This paper presents the results of an experimental investigation on the inception of cavitation in a Quadrant-Edge Orifice meter. The relationship between Reynolds' Number, contraction ratio, and cavitation parameter has been demonstrated under constant back-pressure conditions. Further a graphical method has been suggested for locating exactly the onset of cavitation in a Quadrant-Edge Orifice meter, substantiated with experimental results. This method could as well be applied to other similar pressure differential meters. Values of limiting incipient Reynolds' Numbers are also furnished.

**INTRODUCTION:** Very few investigators have touched upon the effects of cavitation on metering devices and very little information is available regarding the inception of cavitation in such meters. Hence an attempt has been made to study these and this paper mainly deals with the inception of cavitation in Quadrant-Edge Orifice meter.

So far, there is no universally accepted simple method for judging the onset of cavitation. The method that is commonly adopted is visual examination of low pressure zones. This is likely to introduce personal errors. Recently, optical, acoustical and electrical methods have been suggested. These methods generally involve the measurements of the effects of cavitation, which are naturally after inception, and the use of very expensive instruments and hence are not encouraging. Further acoustical methods fail sometimes because of the vibration noise and other noises of the power machines used in the system.

Fig. 1 Photograph of Single Wedge Wings attached to Rocket Booster.





**SPECIFICATION OF THE METERS USED:** The Quadrant-Edge Orifice plates used in the present investigation have the following specifications:

* $\beta$ ratio	$\frac{r}{d}$	Manometer tapping
1. 0.225	0.100	D-D/2 tappings
2. 0.400	0.114	
3. 0.500	0.135	
4. 0.600	0.210	
5. 0.630	0.380	

\*Notations are explained at the end of this paper

**EXPERIMENTAL SET-UP:** The set up is shown in figure 1. A centrifugal pump of capacity 900 gallons per minute at a head of 120 feet was used to supply the water for the system. Two by-passes of 6" dia. each were used to by-pass the surplus water. The test section consisted of 20 ft. upstream section with 4" dia. G.I. pipe and about 25 ft. downstream section with 4" dia. pipe out of which 4 feet was made of transparent perspex sheet and the rest of G.I. The water was collected in a calibrated collecting tank of capacity of 150 cubic feet.

The downstream end was fitted with a gate valve and a mercury-water differential manometer was used to register the pressure at a point 4" upstream of this gate valve.

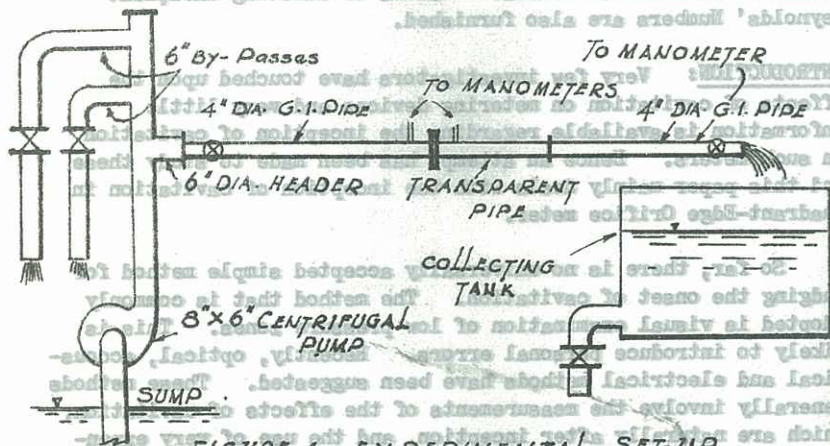


FIGURE 1. EXPERIMENTAL SET-UP

**EXPERIMENTAL PROCEDURE AND DATA COLLECTION:** Pressures  $p_1$  and  $p_2$  were measured for various discharges keeping  $p_3$  as a constant value. This procedure is repeated for various constant  $p_3$  values (Refer Fig. 2).

The onset of cavitation was determined by hearing the characteristic noise as well as by visual observation for checking the validity of the method suggested.

**ANALYSIS OF THE PROBLEM:** The problem is solved under the principle of simple hydrodynamic aspect of the flow system. Referring to Figure 2, the dynamic equilibrium between  $p_1$ ,  $p_2$ ,  $p_3$  and  $V_2$  completely characterises the flow picture across and downstream of the orifice plate. The combination of  $p_1$ ,  $p_2$  and  $V_2$  shows up in the discharge coefficient  $C$  and so in pipe Reynolds' Number  $R_D$ . The downstream conditions can well be represented by  $p_2$  and  $V_2$ , if  $p_3$  is kept as a constant quantity in the analysis. Further in order to represent the cavitation conditions,  $p_2$  and  $V_2$  may be combined with  $p_{sv}$ , the saturated vapour pressure to form a cavitation parameter,  $\sigma$ .

$$\sigma = \frac{p_2 - p_{sv}}{V_2^2 / 2g} \quad (1)$$

Hence a plot of pipe Reynolds' Number against  $\sigma$  completely depicts the flow conditions across and downstream of the orifice plate, provided, the back pressure  $p_3$  is kept constant. Since cavitation entails always in loss of energy, there would be an extra loss of energy in the downstream section under cavitating conditions than under noncavitating, other conditions being similar. For a particular Reynolds' Number, therefore,  $p_2$  will

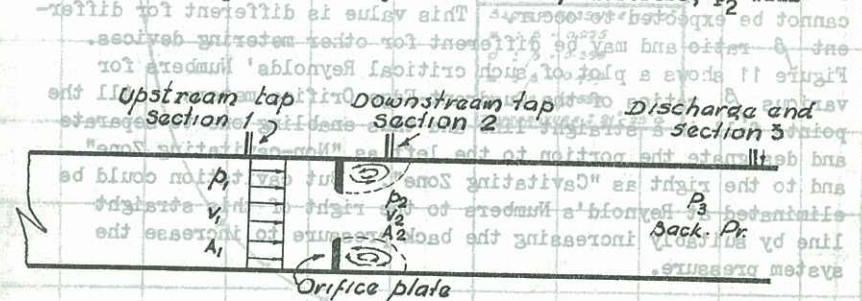


FIGURE 2. DEFINITION SKETCH

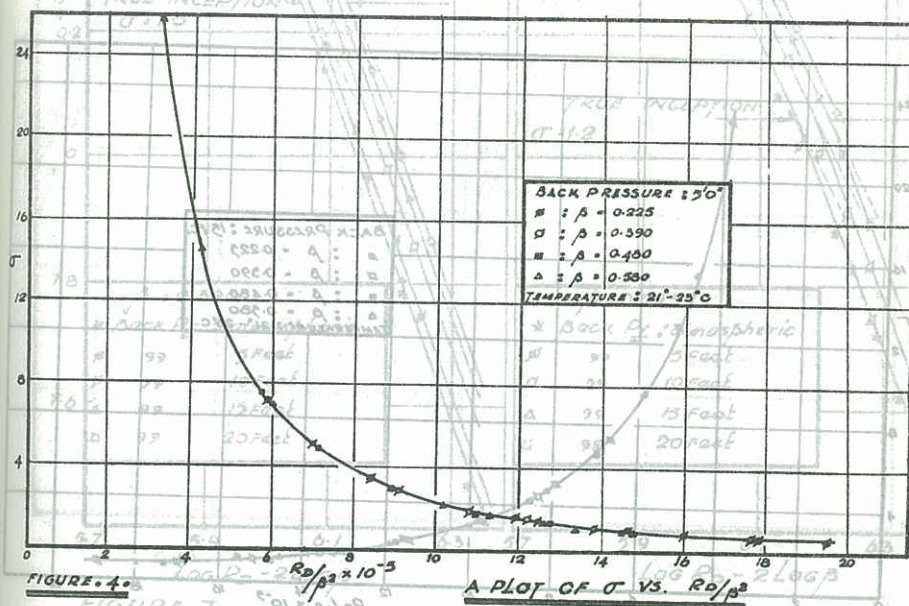
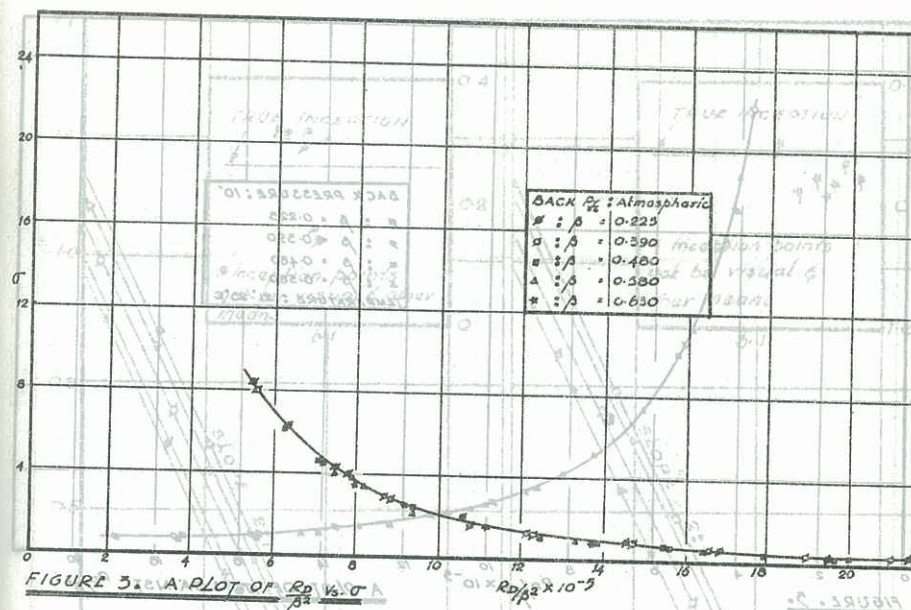


be reduced under cavitating conditions, since  $p_2$  is kept constant. The effect of this reduction in  $p_2$  would be to reduce the value of  $\sigma$ . So one can expect a deviation in the plot of  $\sigma$  versus Reynolds' Number under cavitating conditions. Figs. 3, 4, 5 and 6 show a plot of  $\sigma$  against  $R_D/\beta^2$ , under constant back pressures of 0 (atmospheric), 5, 10 and 15 feet respectively. The division of  $R_D$  by  $\beta^2$  is only to get all the data with a single graph for a constant  $p_2$ . It seems the cavitation parameter holds more or less the same functional relationship with  $R_D$  irrespective of the  $\beta$  ratio, if  $p_2$  is kept constant. Since no significant deviation can be detected from figures 3, 4, 5 and 6, the same data have been plotted in Log-Log scale in figures 7, 8, 9 and 10. It can be seen that the slopes of the straight lines are different under cavitating and noncavitating conditions, with a steeper slope for the latter as per the arguments given above. Obviously the point at which the break occurs represents the conditions at the onset of cavitation. The inception points obtained by visual examination and by hearing the noise are also plotted in the inset block diagrams in each of the figures. It may be noted that these points scatter below the point obtained by the above mentioned graphical solution. This method could be applied to other metering devices as well, since the plots are made to non-pipe dimensional axes.

#### LIMITING REYNOLDS' NUMBER FOR INCEPTION OF CAVITATION.

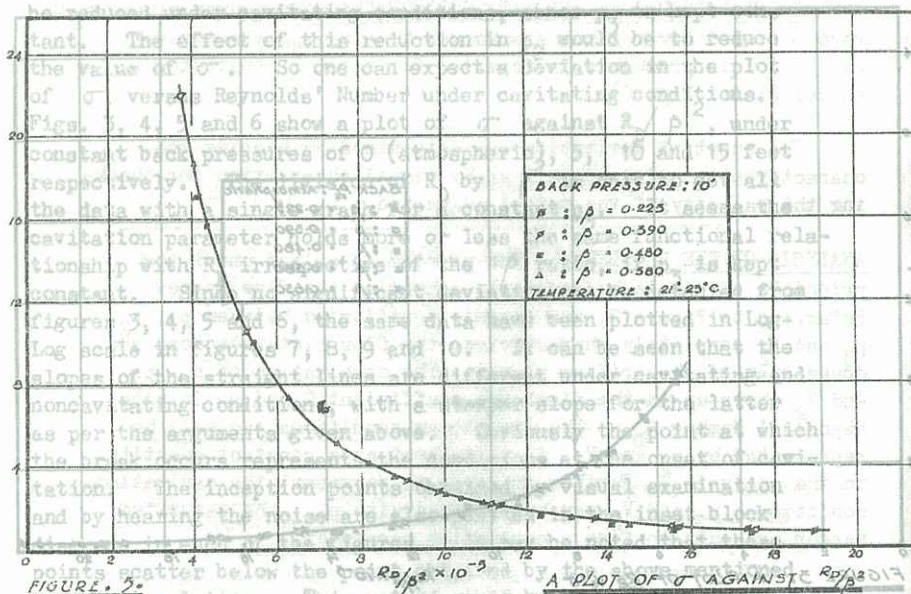
The study has been made for various back-pressures and the minimum back-pressure that may be expected in any metering installation is atmospheric and hence, the Reynolds' Number at which the inception of cavitation takes place with free discharge conditions represents the minimum value below which cavitation cannot be expected to occur. This value is different for different  $\beta$  ratio and may be different for other metering devices. Figure 11 shows a plot of such critical Reynolds' Numbers for various  $\beta$  ratios of the Quadrant-Edge Orifice meters. All the points fall in a straight line and thus enabling one to separate and designate the portion to the left as "Non-cavitating Zone" and to the right as "Cavitating Zone". But cavitation could be eliminated at Reynold's Numbers to the right of this straight line by suitably increasing the back pressure to increase the system pressure.

FIGURE 1. EXPERIMENTAL SET-UP



GRAPHICAL METHOD FOR LOCATING INCEPTION





#### LIMITING REYNOLDS' NUMBER FOR ONSET OF CAVITATION.

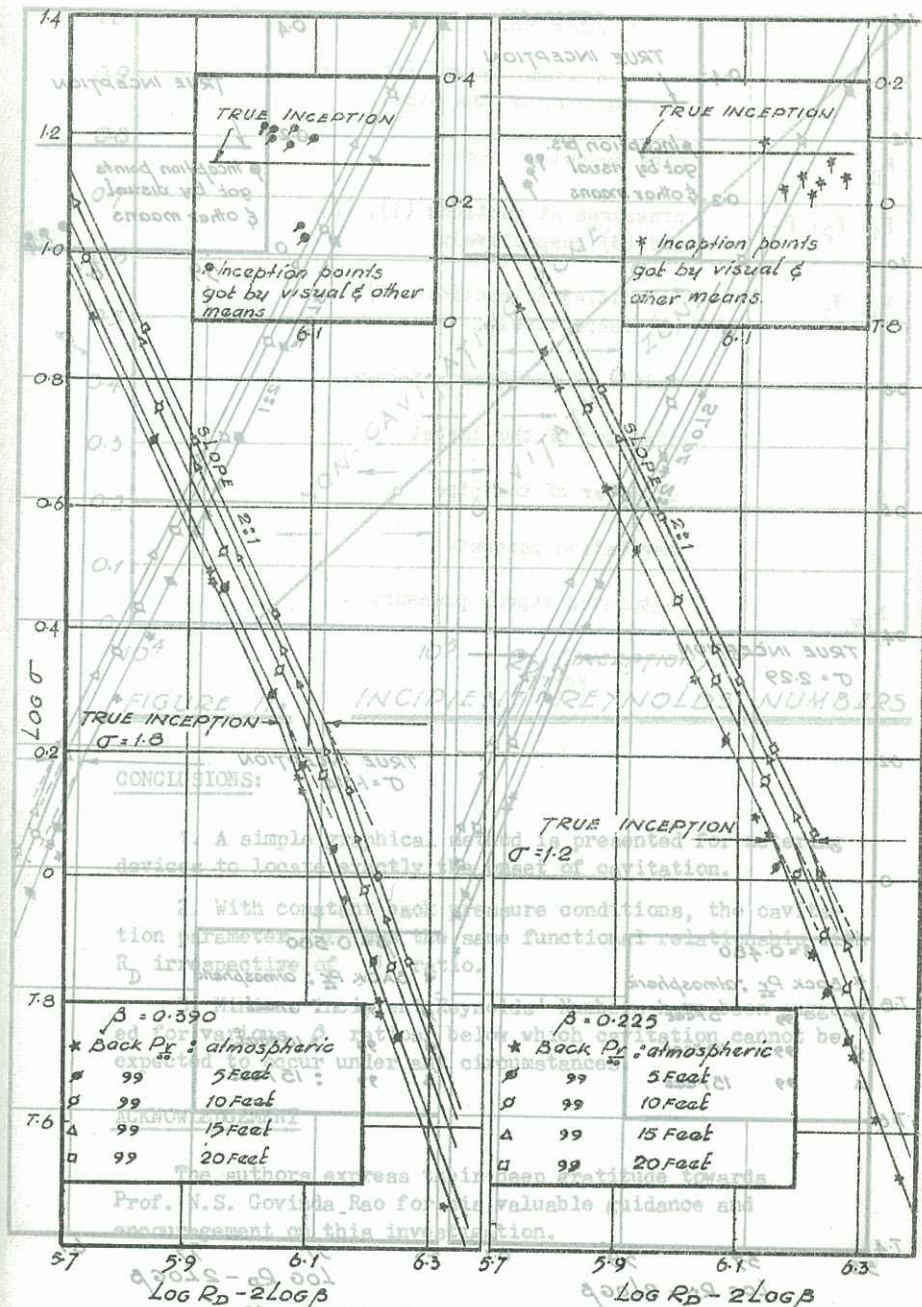
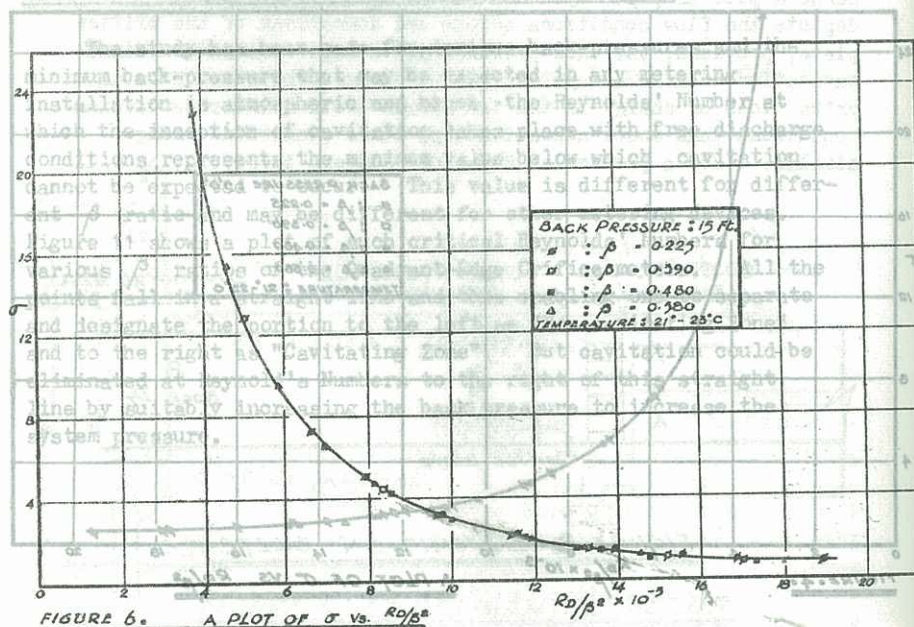


FIGURE 7.

FIGURE 8.

#### GRAPHICAL METHOD FOR LOCATING INCEPTION



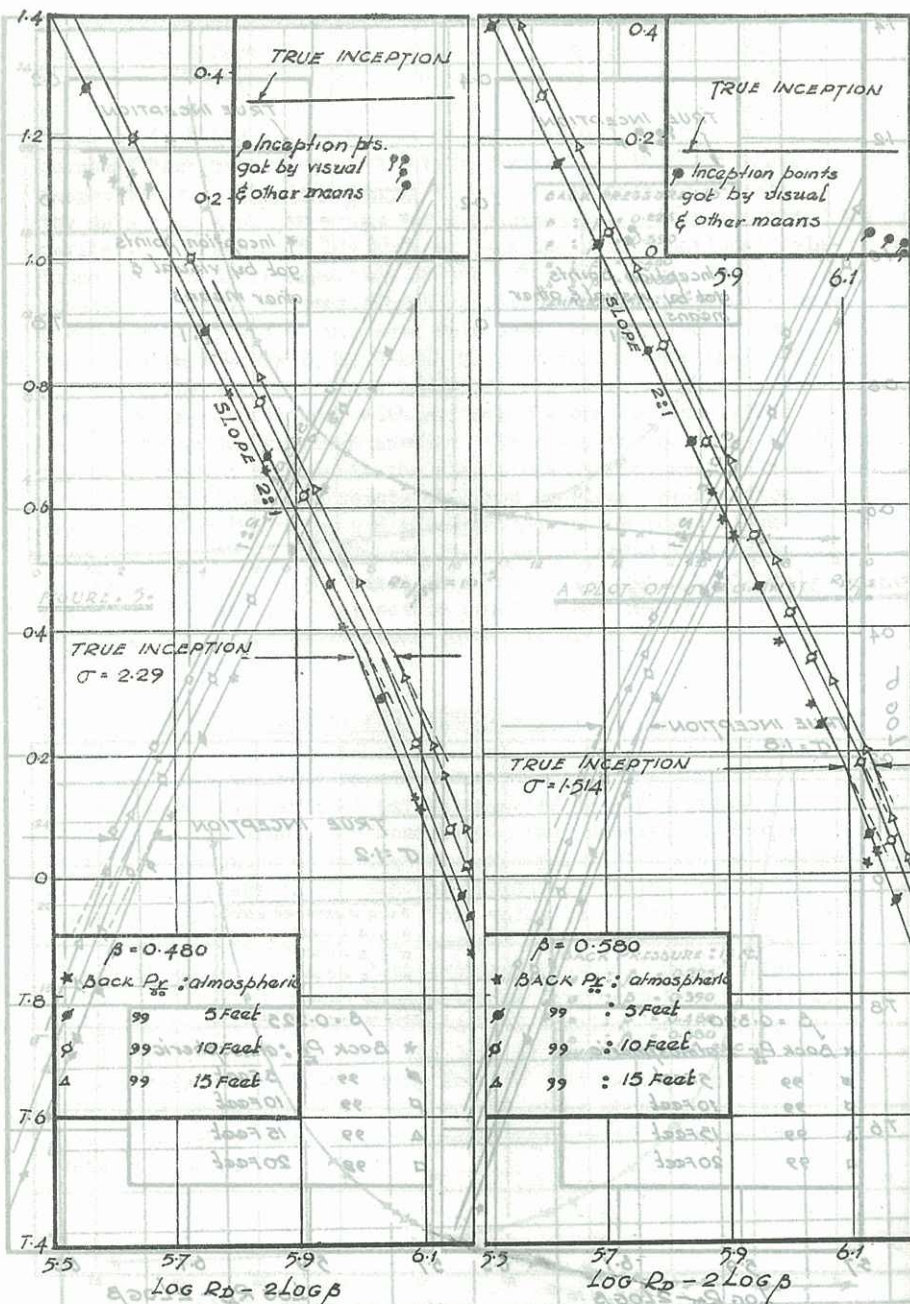


FIGURE 9.

FIGURE 10.

GRAPHICAL METHOD FOR LOCATING INCEPTION.

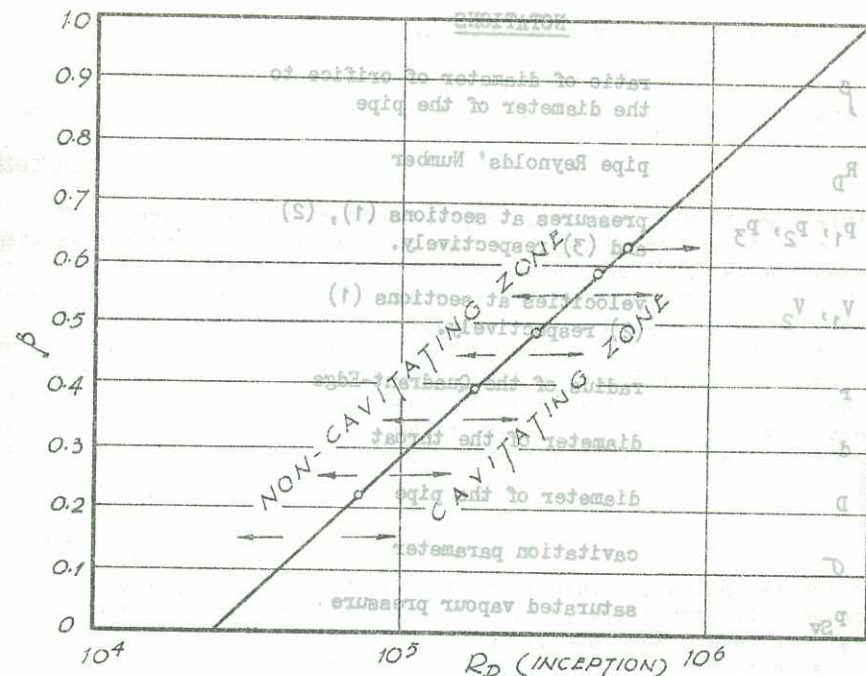


FIGURE 11.

INCIDENT REYNOLDS' NUMBERS

#### CONCLUSIONS:

1. A simple graphical method is presented for metering devices to locate exactly the onset of cavitation.
2. With constant back pressure conditions, the cavitation parameter  $\sigma$  has the same functional relationship with  $R_D$  irrespective of  $\beta$  ratio.
3. Minimum incipient Reynolds' Numbers have been suggested for various  $\beta$  ratios, below which cavitation cannot be expected to occur under any circumstances.

#### ACKNOWLEDGEMENT

The authors express their deep gratitude towards Prof. N.S. Govinda Rao for his valuable guidance and encouragement on this investigation.



## NOTATIONS

$\beta$	ratio of diameter of orifice to the diameter of the pipe
$R_D$	pipe Reynolds' Number
$P_1, P_2, P_3$	pressures at sections (1), (2) and (3) respectively.
$V_1, V_2$	velocities at sections (1) (2) respectively.
$r$	radius of the Quadrant-Edge
$d$	diameter of the throat
$D$	diameter of the pipe
$\sigma$	cavitation parameter
$P_{sv}$	saturated vapour pressure

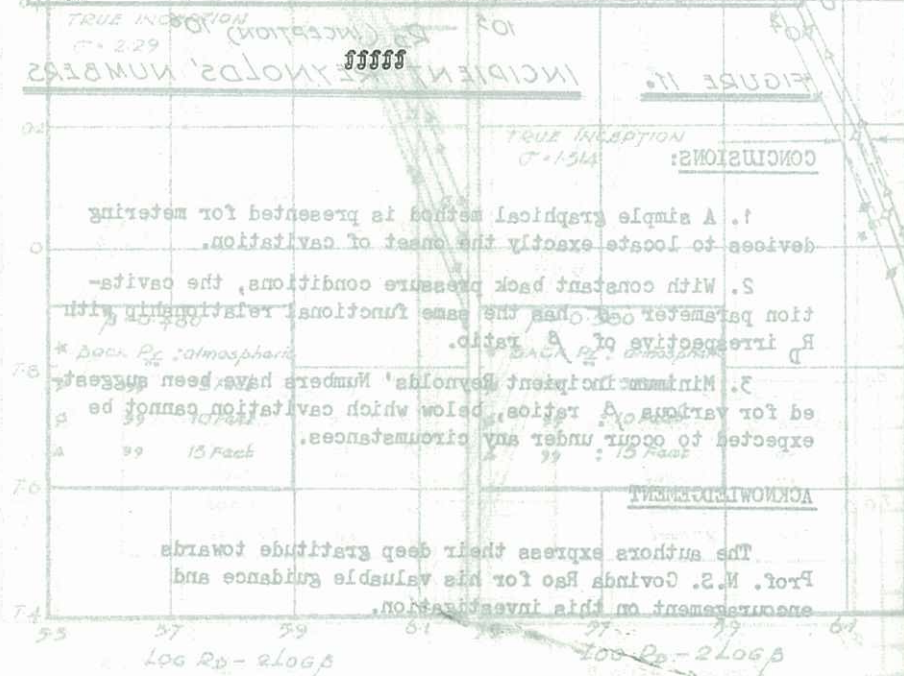


FIGURE 9.

FIGURE 10.

## GRAPHICAL METHOD FOR LOCATING INCEPTION.

## FLOW BELOW A SUBMERGED SLUICE GATE AS A WALL JET PROBLEM

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## Abstract:

This paper presents a study of the submerged flow below a sluice gate as the case of a plane turbulent wall jet under almost zero pressure gradient with a backward flow placed over it. Experiments have been conducted for four supercritical Froude Numbers from 3.01 to 6.44 with the submergence factor varying from 1.60 to 2.24. The velocity distribution in the fully developed region follows closely the curve for the classical wall jet with a slight difference only in the scale factors. The boundary shear stress has been measured with a Preston tube. The flow entrainment and the energy fall have also been considered. The appendix presents some calculations for the classical wall jet.

## Introduction:

The wall jet is defined as a jet of fluid, impinging tangentially (or at an angle) on a boundary, surrounded by stationary (or moving) fluid. The case of the classical wall jet, i.e. the plane turbulent wall jet issuing into the same stationary fluid of semi-infinite extent, on a smooth boundary, is shown in Fig. 1, in which  $y_1$  is the depth of the slot and  $U_1$  is the velocity which is assumed to be uniform for the entire depth of the slot. A very full discussion of the wall jet problem with an account of the other investigations has been given by Rajaratnam (1,2).

For the classical wall jet, the velocity distribution has been found to be fully developed and similar for  $x/y_1 \geq 15$ , where  $x$  is the longitudinal distance from the slot. The length scale at any section is the normal distance from the boundary at which the velocity  $u$  is equal to half the maximum velocity  $u_m$ , and the velocity gradient  $du/dy$  is negative. The velocity scale is the maximum velocity  $u_m$  at that section. This velocity distribution curve has been finalised by Schwarz and Cosart (3) and Rajaratnam (1) using the available experimental information. The nondimensional distance  $y/\delta_1$  is represented by  $\eta$  and the corresponding nondimensional velocity  $u/u_m$  by  $f(\eta)$ .

For the classical wall jet, Sigalla (4) found that

$$\delta_1/y_1 = 0.5 + 0.065 x/y_1 \quad (1)$$

$$u_m/U_1 = 3.45 (x/y_1)^{-0.50} \quad (2)$$

$\tau_0$  is the boundary shear stress at any section, written as

$$\tau_0 = c_f \rho U_m^2/2 \quad (3)$$

where  $c_f$  is the coefficient of skin friction and  $\rho$  is the mass density